

Supplementary Information for

High-resolution solid-state oxygen-17 NMR of actinide-bearing compounds: an insight into the 5*f*-chemistry

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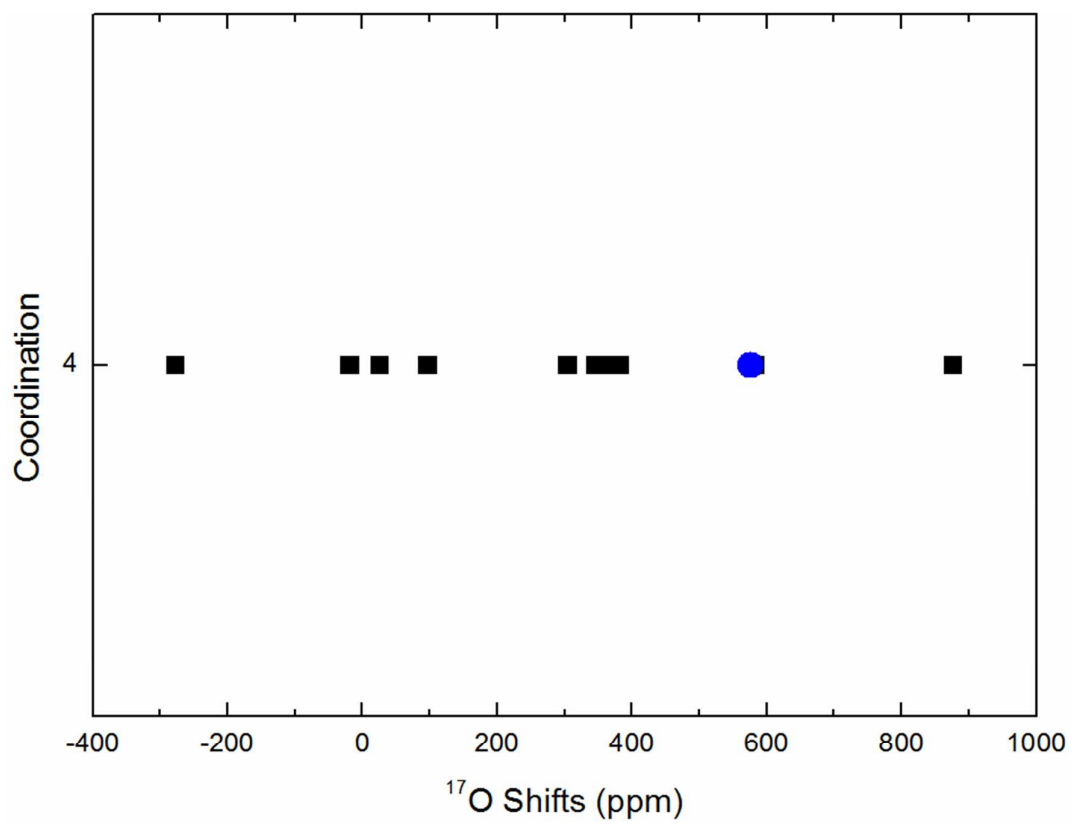


Figure S1: ^{17}O shifts of diamagnetic metal oxides¹ (black squares) compared with that of ThO_2 (blue circle).

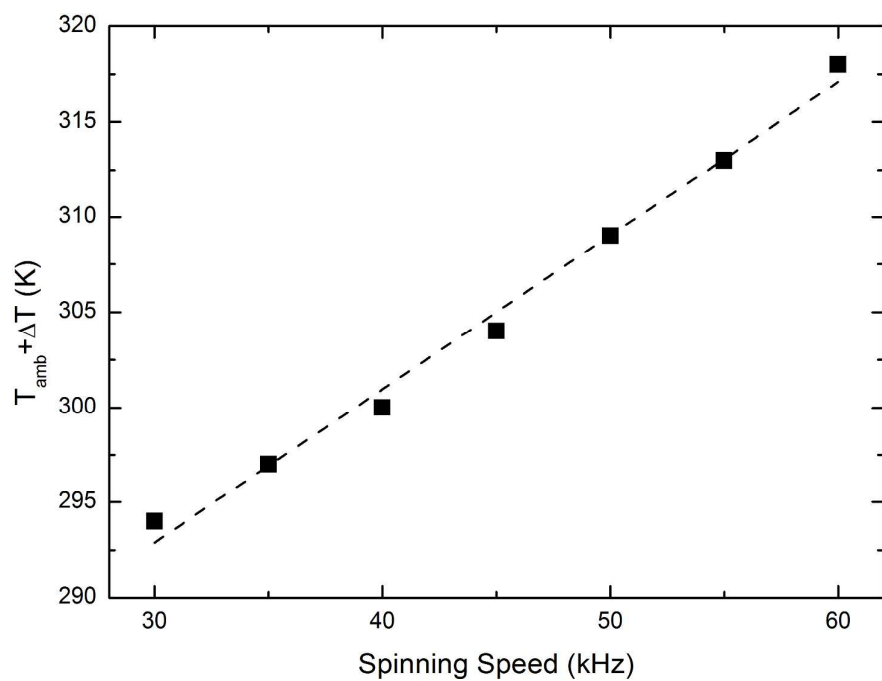


Figure S2: Evolution of the temperature (room temperature (T_{amb}) with a gradient of temperature (ΔT)) with increasing spinning speed. The temperatures have been determined using $\text{Pb}(\text{NO}_3)_2$ correlations². The dashed line is a linear fit.

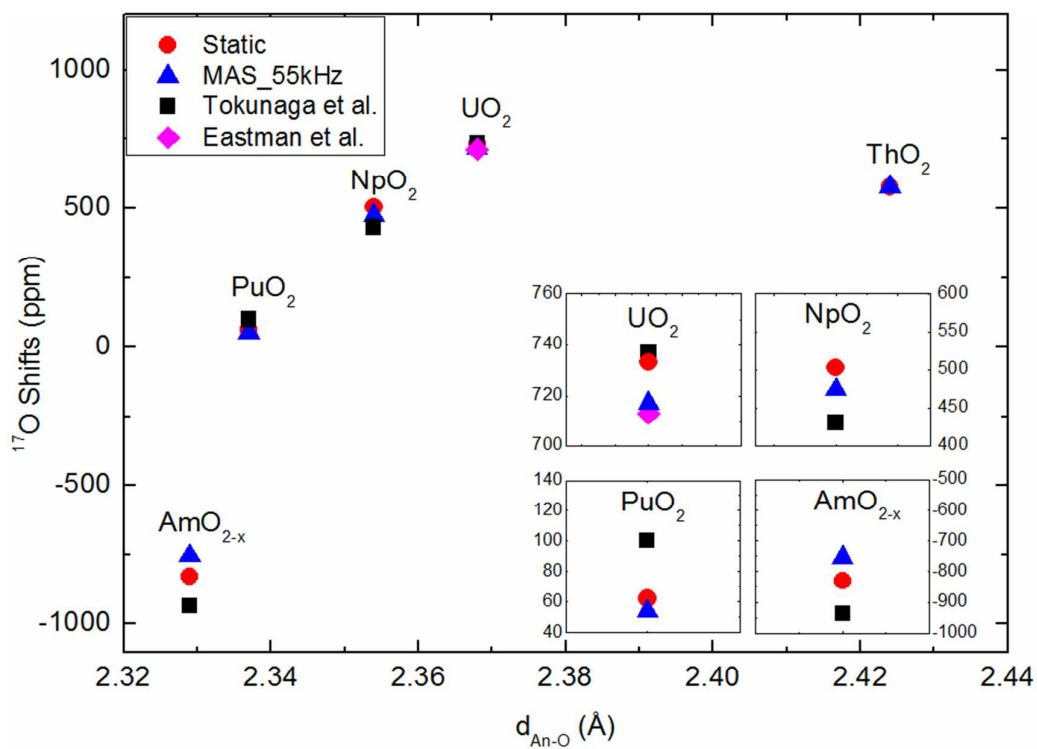


Figure S3: Evolution of the ^{17}O shift as a function of the An-O distance. The results obtained in this study (static and MAS at 55 kHz) are compared with the previous static experiments published by Tokunaga et al.³ (black squares) and Eastman et al. (pink diamond)⁴. The insets give a better view to compare the ^{17}O shifts of each sample obtained in this study and by the other authors.

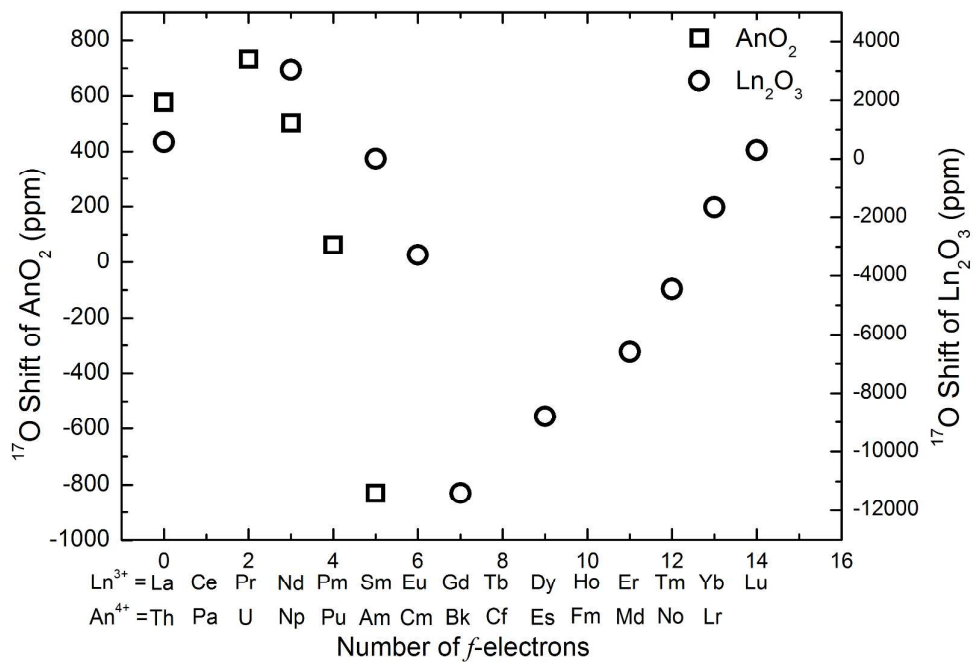


Figure S4: Evolution of the ¹⁷O chemical shifts in lanthanide (Ln₂O₃) (from reference 5) and actinides (AnO₂) dioxides as a function of the number of *f*-electrons.

Table S1: B_4 and B_6 are the crystal field (CF) parameters which appear in H_{CF} (see Equation 4 in main text), and are taken from the given literature references. $A\langle S_z \rangle$ is calculated in several different ways to show the relevant trends: RS = Russell-Saunders coupling, no CF (i.e. only the lowest pure $^{2S+1}L_J$ manifold is accounted for without removing its zero-field degeneracy); IC = Intermediate coupling, no CF (i.e. H_{FI} is diagonalized in order to take into account the mixing of excited $^{2S+1}L'_J$ manifolds into the ground one, which remains $2J+1$ -fold degenerate in zero field); RS+CF = H_{CF} is diagonalized within the lowest pure $^{2S+1}L_J$ manifold, and $A\langle S_z \rangle$ is calculated as a thermal average over the so-obtained wavevectors; IC+CF = $H_{FI} + H_{CF}$ is diagonalized over the whole $5f^n$ configuration states, and $A\langle S_z \rangle$ is calculated as a thermal average over the so-obtained wavevectors. The results of the latter computational scheme are those presented in Figure 5.

| Name | B_4 (cm^{-1}) | B_6 (cm^{-1}) | $A\langle S_z \rangle$ (RS) | $A\langle S_z \rangle$ (RS+CF) | $A\langle S_z \rangle$ (IC) | $A\langle S_z \rangle$ (IC+CF) |
|------------------|-------------------------------|-------------------------------|--------------------------------|-----------------------------------|--------------------------------|-----------------------------------|
| ThO ₂ | | | 0 | 0 | 0 | 0 |
| UO ₂ | -7950 | 3400 | -3.20 | -2.90 | -2.24 | -2.12 |
| NpO ₂ | -8500 | 3400 | -4.91 | -4.42 | -3.13 | -2.97 |
| PuO ₂ | -9750 | 4000 | -4.80 | -4.59 | -1.12 | -1.46 |
| AmO ₂ | -2200 | – | -1.79 | -1.97 | -1.70 | -0.14 |

References

- 1 K. J. D. Mackenzie and M. E. Smith, Multinuclear Solid-State NMR of Inorganic Materials, *Pergamon Materials Series*.
2. Bielecki, A.; Burum, D.P. *J. Magn. Reson. A*, **116**, 215–220 (1995).
3. Tokunaga, Y., et al. NMR studies of actinide dioxides. *J. Alloys Compd.*, **444–445**, 241–245 (2007).
4. Eastman, M. P.; Hecht, H. G.; Lewis, W. B. *J. Chem. Phys.* **54**, 4141-4146 (1971).
5. Yang, S.; Shore, J.; Oldfield, E. *J. Magn. Reson.*, **99**,408-412 (1992).